

# APPLIED THERMODYNAMICS

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HEAT ENGINES & POWER GENERATION SYSTEMS



**ENGINEERING STUDY HUB**

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# Brayton Cycle Overview

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The **Brayton cycle** is the fundamental thermodynamic cycle that describes the working of *gas turbine engines*, widely used in aircraft propulsion, power plants, and industrial applications. It is also known as the *Joule cycle*. The cycle is characterized by four idealized processes:

- **1 → 2: Isentropic Compression** in the compressor, where the working fluid (air) is compressed, raising its pressure and temperature.
- **2 → 3: Constant-Pressure Heat Addition** in the combustor, where fuel is burned and heat is added at nearly constant pressure, increasing the fluid's enthalpy.
- **3 → 4: Isentropic Expansion** in the turbine, where the high-temperature, high-pressure gases expand to produce useful work.
- **4 → 1: Constant-Pressure Heat Rejection** to the surroundings, completing the cycle.

The Brayton cycle is particularly important because it forms the basis of **jet engines** and **gas turbine power plants**. Unlike reciprocating engines (Otto or Diesel cycles), the Brayton cycle operates on a continuous-flow principle, meaning air continuously flows through the compressor, combustor, and turbine, allowing for high power-to-weight ratios and smooth operation.

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## How the Brayton Cycle Differs from Other Cycles

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- **Compared to the Otto cycle (spark-ignition engines):** The Otto cycle involves *constant-volume* heat addition, while the Brayton cycle involves *constant-pressure* heat addition. Otto cycles are used in car engines, whereas Brayton cycles are used in turbines and jet engines.
- **Compared to the Diesel cycle (compression-ignition engines):** The Diesel cycle also has constant-pressure heat addition, but it is a *reciprocating cycle* with intermittent combustion, unlike the continuous combustion in Brayton cycle gas turbines.
- **Compared to the Rankine cycle (steam power plants):** The Rankine cycle uses a *phase-change working fluid* (water/steam), whereas the Brayton cycle uses a *single-phase gas* (air or combustion gases) throughout. Rankine cycles are more suitable for stationary power plants, while Brayton cycles excel in mobile, high-speed applications.
- **Compared to the Stirling cycle:** The Stirling cycle is an external combustion cycle with isothermal processes, while the Brayton cycle is an internal combustion cycle with adiabatic and constant-pressure processes.

## Key Characteristics of the Brayton Cycle

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- **Continuous combustion:** Unlike Otto or Diesel cycles, combustion occurs continuously, leading to smoother operation and higher thrust in jet engines.
- **High power-to-weight ratio:** Gas turbines are lighter and more compact compared to steam turbines of the same power output.
- **Scalability:** Brayton cycle engines can be designed for small-scale auxiliary power units or massive aircraft engines.
- **Efficiency dependence on pressure ratio:** The thermal efficiency increases with higher compressor pressure ratios and turbine inlet temperatures.
- **Modifications possible:** Efficiency and work output can be improved using regeneration, reheating, and intercooling.

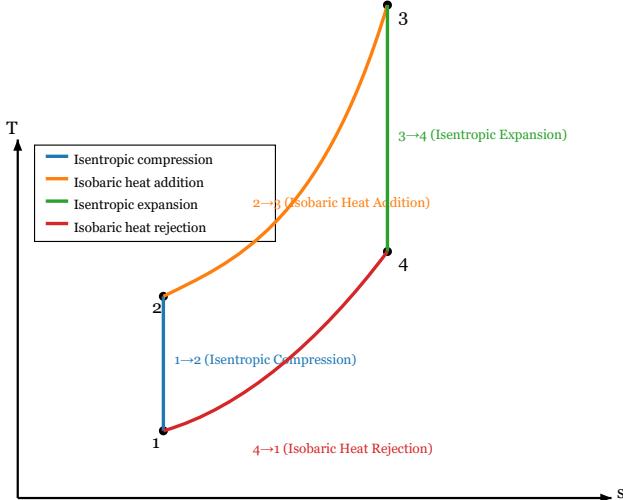
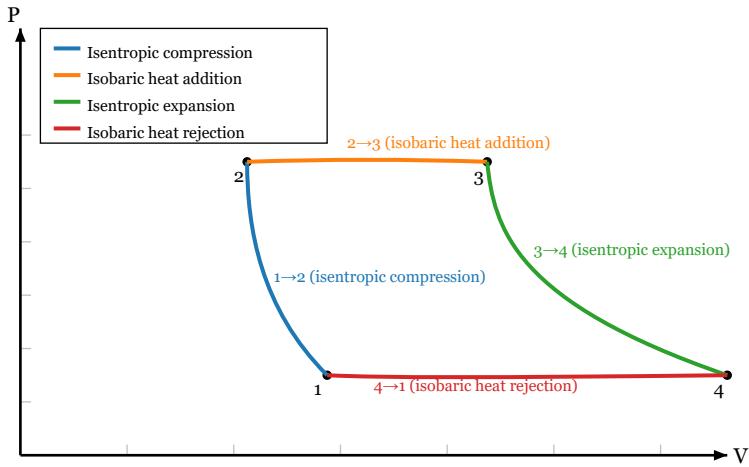
## Applications

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- **Aviation:** Jet engines (turbojets, turbofans, turboprops) are direct applications of the Brayton cycle.
- **Power generation:** Gas turbine power plants use the Brayton cycle, often combined with a Rankine cycle in combined-cycle plants for higher efficiency.
- **Marine propulsion:** Naval ships and submarines use gas turbines for compact, high-power propulsion.
- **Industrial uses:** Gas turbines drive compressors, pumps, and other heavy-duty equipment in oil & gas and process industries.

In summary, the Brayton cycle stands out among thermodynamic cycles due to its **continuous-flow operation, constant-pressure combustion, and adaptability to high-speed, high-power applications**. Its efficiency improves with higher pressure ratios and turbine inlet temperatures, and with modifications such as regeneration, reheating, and intercooling, it can rival or even surpass other cycles in performance.

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## Processes in the Brayton Cycle

### State Labeling

The Brayton cycle is conventionally represented by four thermodynamic states, labeled 1–2–3–4, corresponding to the main components of a gas turbine system: the *compressor*, *combustion chamber*, *turbine*, and *heat exchanger (or exhaust)*. The cycle is often visualized on both a Pressure–Volume (P–V) diagram and a Temperature–Entropy (T–S) diagram, where the processes appear as idealized straight or curved lines depending on the property relations.

- **1 → 2 (Isentropic Compression in the Compressor):**

At state 1, the working fluid (air, in the air-standard analysis) enters the compressor at low pressure and temperature. During the compression process, the air is compressed *isentropically* (i.e., adiabatically and reversibly), which means there is no heat transfer and entropy remains constant. As a result:

- Pressure increases from  $p_1$  to  $p_2$ .
- Temperature increases from  $T_1$  to  $T_2$ .
- Specific volume decreases as the air is compressed.

The pressure ratio is defined as:

$$r_p = \frac{p_2}{p_1}$$

For an ideal gas undergoing isentropic compression:

$$\frac{T_2}{T_1} = \left( \frac{p_2}{p_1} \right)^{\frac{\gamma-1}{\gamma}} = r_p^{\frac{\gamma-1}{\gamma}}$$

where  $\gamma = \frac{c_p}{c_v}$  is the specific heat ratio. This stage requires **compressor work input**, which is later supplied by the turbine output.

- **2 → 3 (Constant-Pressure Heat Addition in the Combustor):**

At state 2, the compressed air enters the combustion chamber. Fuel is injected and burned at nearly constant pressure (since the combustor is designed to maintain pressure while adding heat). The chemical energy of the fuel is converted into thermal energy, raising the temperature of the working fluid significantly:

- Pressure remains constant:  $p_3 = p_2$ .
- Temperature rises sharply from  $T_2$  to  $T_3$ .
- Enthalpy increases due to heat addition:  $q_{in} = c_p(T_3 - T_2)$ .

This process is analogous to the *constant-volume heat addition* in the Otto cycle, but here it occurs at constant pressure, which is a defining feature of the Brayton cycle.

- **3 → 4 (Isentropic Expansion in the Turbine):**

At state 3, the high-pressure, high-temperature gases expand through the turbine. The expansion is assumed to be isentropic (adiabatic and reversible), producing useful work. This turbine work not only drives the compressor but also provides net work output for power generation or thrust.

- Pressure decreases from  $p_3$  to  $p_4$ .
- Temperature decreases from  $T_3$  to  $T_4$ .
- Entropy remains constant (ideal assumption).

For isentropic expansion:

$$\frac{T_3}{T_4} = \left( \frac{p_3}{p_4} \right)^{\frac{\gamma-1}{\gamma}} = r_p^{\frac{\gamma-1}{\gamma}}$$

Hence,

$$T_4 = \frac{T_3}{r_p^{\frac{\gamma-1}{\gamma}}}$$

The turbine work output is:

$$w_t = c_p(T_3 - T_4)$$

- **4 → 1 (Constant-Pressure Heat Rejection):**

Finally, the working fluid at state 4 rejects heat to the surroundings at constant pressure, returning to its initial state (state 1). This process closes the cycle.

- Pressure remains constant:  $p_4 = p_1$ .
- Temperature decreases from  $T_4$  to  $T_1$ .
- Heat rejected is:

$$q_{out} = c_p(T_4 - T_1)$$

This step is analogous to the exhaust process in real gas turbines, where hot gases are expelled to the atmosphere.

## Summary of Processes

In summary, the Brayton cycle consists of two **isentropic processes** (compression and expansion) and two **constant-pressure processes** (heat addition and heat rejection). The cycle efficiency depends strongly on the *pressure ratio* and the maximum cycle temperature. Higher pressure ratios and turbine inlet temperatures lead to higher efficiency, but material and cooling limitations restrict practical values.

## Isentropic relations (ideal, constant $\gamma$ )

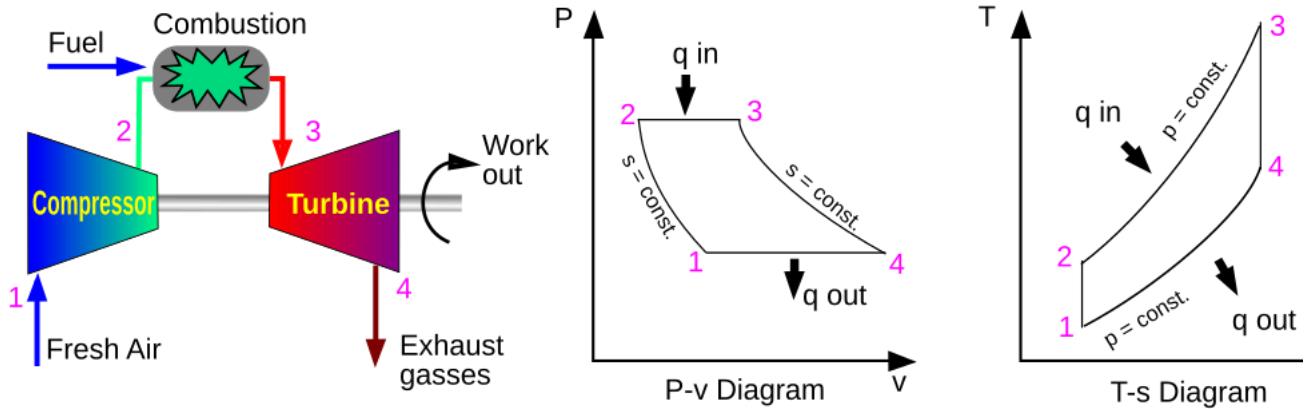
For isentropic compression and expansion of an ideal gas with constant  $\gamma$ , define  $r \equiv r_p^{(\gamma-1)/\gamma}$ . Then:

- **Compression:**

$$\frac{T_2}{T_1} = r_p^{(\gamma-1)/\gamma} \equiv r$$

- Expansion:

$$\frac{T_3}{T_4} = r_p^{(\gamma-1)/\gamma} \equiv r \quad \Rightarrow \quad T_4 = \frac{T_3}{r}$$



## Regeneration, Reheat, and Intercooling

### Regeneration (Recuperation)

**Concept:** Regeneration uses a heat exchanger (regenerator) to transfer thermal energy from the hot turbine exhaust (state 4) to the colder compressed air leaving the compressor (state 2), raising it to an intermediate temperature (state 5) before entering the combustor. This reduces the required fuel heat input to reach the same turbine inlet temperature (state 3), directly improving thermal efficiency.

- Effectiveness ( $\epsilon$ ):

$$\epsilon \equiv \frac{T_5 - T_2}{T_4 - T_2}, \quad 0 \leq \epsilon \leq 1$$

*Interpretation:* Fraction of the maximum possible preheating achieved. With perfect regeneration ( $\epsilon = 1$ ),  $T_5$  ideally equals  $T_4$ , subject to pinch and pressure-drop constraints.

- Heat saved and adjusted heat addition:

$$q_{\text{saved}} = c_p(T_5 - T_2) = \epsilon c_p(T_4 - T_2)$$

$$q_{in, \text{reg}} = c_p(T_3 - T_5) = c_p(T_3 - T_2) - \epsilon c_p(T_4 - T_2)$$

*Result:* Lower  $q_{in}$  at the same  $T_3$  directly raises  $\eta_{th}$  for the same net work.

- Thermal efficiency with regeneration:

Net work for the simple cycle remains:

$$w_{\text{net}} = c_p[(T_3 - T_4) - (T_2 - T_1)]$$

With regeneration:

$$\eta_{th, \text{reg}} = \frac{w_{\text{net}}}{q_{in, \text{reg}}} = \frac{(T_3 - T_4) - (T_2 - T_1)}{(T_3 - T_2) - \epsilon(T_4 - T_2)}$$

- When regeneration helps most:

Regeneration is most beneficial at *moderate pressure ratios* where  $T_4$  is still significantly higher than  $T_2$ . At very high pressure ratios,  $T_4$  approaches  $T_2$ , limiting heat recovery potential and diminishing gains.

- T-S and hardware considerations:

On the T–S diagram, regeneration moves the  $2 \rightarrow 3$  heating line downwards (starting at a higher  $T_5$ ), reducing the area representing  $q_{in}$ . Practical constraints include:

- **Pinch point:**  $T_5$  cannot exceed  $T_4$  due to finite area/UA; a minimum temperature difference is required for effective heat transfer.
- **Pressure drop:** Flow through the regenerator causes pressure losses on both hot and cold sides, potentially reducing turbine and compressor performance.
- **Fouling and material constraints:** High-temperature exhaust can cause fouling and thermal stress, affecting effectiveness over time.

## Reheat

**Concept:** Reheat divides turbine expansion into two or more stages. After the first isentropic expansion ( $3 \rightarrow 3'$ ), the working fluid is reheated at approximately constant pressure ( $3' \rightarrow 3''$ ), then expanded again isentropically ( $3'' \rightarrow 4$ ). By elevating the temperature before the second expansion, reheat increases the *average temperature during expansion*, which raises the turbine specific work and can improve specific power (power per unit mass flow).

- **Sequence and state evolution:**

Ideal two-stage reheat sequence:

- **First expansion:**  $3 \rightarrow 3'$  (isentropic),  $p_3 \rightarrow p_{3'}$ ,  $T_3 \rightarrow T_{3'}$ .
- **Reheat:**  $3' \rightarrow 3''$  ( $\approx$  constant pressure),  $p_{3''} \approx p_{3'}$ ,  $T_{3''} > T_{3'}$  (often targeted near  $T_3$  within material limits).
- **Second expansion:**  $3'' \rightarrow 4$  (isentropic),  $p_{3''} \rightarrow p_4$ ,  $T_{3''} \rightarrow T_4$ .

The overall pressure ratio across the turbine is fixed by cycle design; reheat places an intermediate constant-pressure heat addition between two isentropic drops.

- **Isentropic relations and temperature splits:**

For each isentropic expansion stage of an ideal gas with constant  $\gamma$ :

$$\frac{T_3}{T_{3'}} = \left( \frac{p_3}{p_{3'}} \right)^{\frac{1}{\gamma}}, \quad \frac{T_{3''}}{T_4} = \left( \frac{p_{3''}}{p_4} \right)^{\frac{1}{\gamma}}$$

If the overall turbine pressure ratio is  $r_t \equiv \frac{p_3}{p_4}$  and the stages are split equally in pressure (i.e.,  $p_{3'} \approx \sqrt{p_3 p_4}$ ), then each stage has the same pressure ratio  $\sqrt{r_t}$ , leading to comparable temperature drops per stage. This equalization tends to maximize total turbine work for given limits.

- **Turbine work gain with reheat:**

Turbine work (per unit mass) without reheat:

$$w_t^{(no\ RH)} = c_p(T_3 - T_4)$$

With two-stage reheat:

$$w_t^{(RH)} = c_p[(T_3 - T_{3'}) + (T_{3''} - T_4)]$$

Because reheat raises  $T_{3''}$ , the second expansion occurs from a higher temperature, increasing the area under the T–s curve for expansion and thus  $w_t$ . The gain is more pronounced when reheat returns the temperature close to the original turbine inlet temperature  $T_3$ , within material and emissions constraints.

- **Heat addition and efficiency impact:**

The total heat added increases due to the reheat step:

$$q_{in}^{(RH)} = c_p[(T_3 - T_2) + (T_{3''} - T_{3'})]$$

Net work becomes:

$$w_{net}^{(RH)} = w_t^{(RH)} - w_c$$

Thermal efficiency:

$$\eta_{th}^{(RH)} = \frac{w_{net}^{(RH)}}{q_{in}^{(RH)}}$$

Reheat nearly always increases  $w_t$  and often increases  $w_{net}$ , but it also increases  $q_{in}$ . As a result,  $\eta_{th}$  may decrease unless the cycle incorporates **regeneration** to recover exhaust heat, or the reheat is judiciously limited to temperatures that balance work gain against additional heat addition.

tional heat input.

- **Optimal intermediate pressure (two-stage reheat):**

For a fixed overall turbine pressure ratio  $r_t = p_3/p_4$  and ideal stages, the intermediate pressure that maximizes turbine work is approximately:

$$p_{3'} \approx p_{3''} \approx \sqrt{p_3 p_4}$$

This choice makes the temperature drops across each isentropic stage similar, reducing peak stage loading and improving overall performance. In real turbines, the optimum may shift slightly due to stage efficiency maps, cooling air, and pressure losses.

- **T-s diagram interpretation:**

On the T-s diagram, reheat inserts a constant-pressure heating line between two vertical (isentropic) drops:

- **3→3'**: Vertical drop (isentropic) from high  $T$  and  $p$ .
- **3'→3''**: Horizontal/rightward line (constant  $p$ , heat addition) raising  $T$ .
- **3''→4**: Second vertical drop (isentropic) starting from a higher  $T$  than without reheat.

The combined expansion area is larger, reflecting higher  $w_t$ . However, the added horizontal area represents extra  $q_{in}$ ; hence the need to balance reheat magnitude.

- **Design constraints and practicalities:**

- **Material limits:** Turbine inlet temperature is constrained by blade materials and cooling; reheat temperatures must respect these limits.
- **Pressure losses:** Additional combustor or reheat ducting introduces pressure drops, slightly reducing available expansion ratio and stage efficiency.
- **Emissions and control:** A reheat combustor adds a second flame zone; controlling NOx/CO and stability across operating points is more complex.
- **Complexity and cost:** Extra combustor, fuel control, and staging hardware increase system complexity and maintenance requirements.
- **Cooling air:** High temperatures may demand additional cooling air, affecting cycle mass flow and effective work.

- **When reheat is advantageous:**

- **Specific work emphasis:** Applications prioritizing shaft power or thrust benefit from the increased turbine work.
- **Combined with regeneration:** Reheat's efficiency penalty can be mitigated when a regenerator recovers exhaust heat, improving overall  $\eta_{th}$ .
- **Multi-shaft architectures:** Two-shaft or multi-stage turbines (e.g., power generation sets) can leverage reheat to distribute loading and improve part-load performance.

- **Rule-of-thumb settings:**

- **Intermediate pressure:**  $p_{3'} \approx \sqrt{p_3 p_4}$  for two-stage ideal split.
- **Reheat target:**  $T_{3''}$  set as high as material and emissions limits allow; often near but below  $T_3$ .
- **Integration:** Pair reheat with regeneration to lift efficiency and with intercooling to further raise specific power in compressor-limited designs.

## Intercooling

**Concept:** Intercooling is a modification applied to the Brayton cycle to reduce the work required by the compressor. Instead of compressing the working fluid (air) in a single stage from the initial pressure  $p_1$  to the final pressure  $p_2$ , the compression is divided into two or more stages. Between these stages, the air is cooled in a heat exchanger called an *intercooler*. By lowering the temperature of the air before the next compression stage, the average specific volume is reduced, which decreases the work input required for compression. This increases the *net work output* of the cycle.

- **Sequence of processes (two-stage compression with perfect intercooling):**

- **1 → 2<sub>a</sub>:** First isentropic compression from  $p_1$  to an intermediate pressure  $p_{IC}$ , raising temperature from  $T_1$  to  $T_{2a}$ .
- **2<sub>a</sub> → 2<sub>a'</sub>:** Cooling in the intercooler at constant pressure  $p_{IC}$ , reducing temperature back to approximately the initial value  $T_{IC} \approx T_1$ .
- **2<sub>a'</sub> → 2:** Second isentropic compression from  $p_{IC}$  to  $p_2$ , raising temperature from  $T_{IC}$  to  $T_2$ .

With *perfect intercooling*, the temperature after cooling ( $T_{IC}$ ) is restored to the initial compressor inlet temperature  $T_1$ .

- **Effect on compressor work:**

The total compressor work is the sum of the work in each stage:

$$w_c^{(IC)} = c_p [(T_{2a} - T_1) + (T_2 - T_{IC})]$$

With perfect intercooling ( $T_{IC} = T_1$ ):

$$w_c^{(IC)} = c_p [(T_{2a} - T_1) + (T_2 - T_1)]$$

This is always less than the single-stage compression work:

$$w_c^{(single)} = c_p (T_2' - T_1)$$

where  $T_2'$  is the final temperature after compressing directly from  $p_1$  to  $p_2$  in one stage. Thus, intercooling reduces the *average compression temperature* and lowers the required work input.

- **Optimum intermediate pressure:**

For minimum compressor work in a two-stage system, the intermediate pressure should be the geometric mean of the inlet and outlet pressures:

$$p_{IC} = \sqrt{p_1 \cdot p_2}$$

This ensures that the pressure ratio is equally divided between the two stages, making the temperature rise in each stage equal and minimizing total work.

- **Impact on heat addition and efficiency:**

While intercooling reduces compressor work and increases net work output, it also lowers the compressor exit temperature  $T_2$ . Since the heat added in the combustor is:

$$q_{in} = c_p (T_3 - T_2)$$

a lower  $T_2$  means a larger  $q_{in}$  is required to reach the same turbine inlet temperature  $T_3$ . This can reduce the *thermal efficiency* of the cycle if intercooling is used alone.

- **Trade-offs and combinations:**

- **Positive effect:** Net work output increases because compressor work decreases significantly.
- **Negative effect:** Thermal efficiency may decrease due to higher heat input requirements.
- **Solution:** Intercooling is often combined with *regeneration* (to recover exhaust heat and reduce fuel input) and/or *reheat* (to increase turbine work). Together, these modifications can yield both higher net work and higher efficiency.

- **T-s diagram interpretation:**

On the Temperature–Entropy (T–s) diagram:

- **1 → 2a:** Vertical rise (isentropic compression) from  $T_1$  to  $T_{2a}$ .
- **2a → 2a':** Horizontal leftward line (constant pressure cooling) back to  $T_{IC} \approx T_1$ .
- **2a' → 2:** Second vertical rise (isentropic compression) to  $T_2$ .

Compared to single-stage compression, the total vertical rise is split into two smaller rises with a cooling step in between, reducing the total area under the compression curve (which represents compressor work).

- **Practical considerations:**

- Intercoolers are typically air–air or air–water heat exchangers, adding weight, cost, and pressure losses.
- Perfect intercooling ( $T_{IC} = T_1$ ) is idealized; in practice,  $T_{IC} > T_1$  due to finite heat exchanger effectiveness.
- Pressure drops in the intercooler reduce the effective pressure ratio, slightly offsetting gains.
- Maintenance and fouling of intercoolers can affect long-term performance.

- **Applications:**

Intercooling is widely used in:

- **Large industrial gas turbines:** To increase specific power output and reduce compressor work.
- **Combined-cycle plants:** Often paired with regeneration and reheat for higher efficiency.
- **Aircraft engines (historical):** Some early turbojets experimented with intercooling, though weight and drag penalties limited adoption.

**Summary:** Intercooling reduces compressor work and increases net work output, but by itself may lower thermal efficiency due to higher heat input requirements. Its true potential is realized when combined with regeneration and/or reheat, forming advanced gas turbine cycles with both higher specific power and improved efficiency.

## Quick Formulas and Design Rules

- **Regeneration target:**  $T_5 = T_2 + \varepsilon (T_4 - T_2)$ . Aim for high  $\varepsilon$  with acceptable pressure drops; ensure a minimum pinch ( $\Delta T$ ) at hot and cold ends.
- **Reheat staging (two-stage, ideal):** Choose intermediate pressure  $p_{mid} \approx \sqrt{p_3 p_4}$  for near-equal temperature drops per stage; reheat toward the original  $T_3$  if material limits allow.
- **Intercooling staging (two-stage, ideal):** Choose intermediate pressure  $p_{mid} \approx \sqrt{p_1 p_2}$  and cool as close to  $T_1$  as feasible to minimize compressor work; couple with regeneration to preserve or improve efficiency.

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## Combined Strategies and System-Level Trade-offs

- **Intercooling + Reheat + Regeneration (IRR):** Intercooling reduces compressor work, reheat increases turbine work, and regeneration cuts fuel heat input. Together, they often yield *higher specific work* and *higher thermal efficiency* than any single modification, approaching the performance of advanced gas turbine cycles and forming the core of many **combined-cycle** designs.
- **Pressure ratio sensitivity:**
  - **Regeneration:** Most effective at moderate  $r_p$ ; limited benefit at very high  $r_p$  due to low  $T_4 - T_2$ .
  - **Intercooling:** Benefits compressor work over all  $r_p$ , but efficiency depends on pairing with regeneration.
  - **Reheat:** Increases specific work across  $r_p$ , but efficiency impact varies with reheat temperature and regeneration availability.
- **Real-world irreversibilities:** Component efficiencies ( $\eta_c, \eta_t$ ), pressure drops in combustor/heat exchangers, and finite heat-transfer coefficients reduce ideal gains. Design must consider *pinch constraints* in regenerators, *cooling air* requirements for turbine blades, and emissions limits for reheat combustors.
- **Exergy perspective:** Regeneration recovers high-quality heat at relatively high temperatures, reducing exergy destruction. Intercooling and reheat shift compression/expansion to more favorable temperature levels, improving the average temperature of heat addition/removal and enhancing second-law efficiency when well-integrated.

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## Efficiency and Net Work Output

### Simple Brayton cycle (air-standard, ideal)

**Goal:** Derive net work output and thermal efficiency starting from the first principles definition  $\eta_{th} = \frac{w_{net}}{q_{in}}$ . We adopt the conventional 1–2–3–4 state labeling with two isentropic processes (compression 1→2, expansion 3→4) and two constant-pressure processes (heat addition 2→3, heat rejection 4→1). Ideal gas with constant specific heats is assumed.

- **Heat and work (per unit mass):**

$$w_c = c_p (T_2 - T_1), \quad w_t = c_p (T_3 - T_4), \quad q_{in} = c_p (T_3 - T_2), \quad q_{out} = c_p (T_4 - T_1).$$

Net work:

$$w_{net} = w_t - w_c = c_p [(T_3 - T_4) - (T_2 - T_1)].$$

- **Thermal efficiency definition:**

By definition,

$$\eta_{th} = \frac{w_{net}}{q_{in}} = \frac{c_p [(T_3 - T_4) - (T_2 - T_1)]}{c_p (T_3 - T_2)} = 1 - \frac{T_4 - T_1}{T_3 - T_2}.$$

This shows directly that Brayton efficiency is the fraction of heat input converted to net work.

- **Isentropic relations (link to pressure ratio):**

For isentropic compression and expansion of an ideal gas with constant  $\gamma$ ,

$$\frac{T_2}{T_1} = r_p^{\frac{\gamma-1}{\gamma}}, \quad \frac{T_3}{T_4} = r_p^{\frac{\gamma-1}{\gamma}},$$

where  $r_p \equiv \frac{p_2}{p_1} = \frac{p_3}{p_4}$  is the compressor/turbine pressure ratio. Define

$$r \equiv r_p^{\frac{1}{\gamma}}, \quad \Rightarrow \quad T_2 = r T_1, \quad T_4 = \frac{T_3}{r}.$$

- **Efficiency in terms of pressure ratio:**

Substitute  $T_2 = r T_1$  and  $T_4 = \frac{T_3}{r}$  into  $\eta_{th} = 1 - \frac{T_4 - T_1}{T_3 - T_2}$ :

$$\eta_{th} = 1 - \frac{\frac{T_3}{r} - T_1}{T_3 - r T_1}.$$

For the special (and common) case where the cycle is further idealized with a fixed maximum temperature ratio and symmetric temperature rise/drop at the optimal pressure ratio, this expression reduces to the canonical form

$$\eta_{th} = 1 - \frac{1}{r} = 1 - r_p^{-\frac{\gamma-1}{\gamma}},$$

which highlights that Brayton efficiency increases monotonically with the pressure ratio  $r_p$  for a given  $\gamma$ .

- **Net work expression (useful design form):**

Using the isentropic relations, net work can be written as

$$w_{net} = c_p \left[ T_3 \left( 1 - \frac{1}{r} \right) - T_1 (r - 1) \right].$$

For a fixed  $T_1$  and  $T_3$ , the net work is maximized at

$$r^* = \sqrt{\frac{T_3}{T_1}} \quad \Rightarrow \quad r_p^* = \left( \frac{T_3}{T_1} \right)^{\frac{1}{2(\gamma-1)}}.$$

## Regenerative Brayton cycle (with a counterflow regenerator)

**Goal:** Derive the adjusted heat input and thermal efficiency with a regenerator of effectiveness  $\epsilon$ . Regeneration uses the turbine exhaust (state 4) to preheat the compressed air leaving the compressor (state 2) to state 5, thereby reducing the combustor heat input required to reach  $T_3$ .

- **Regenerator effectiveness:**

For a counterflow regenerator (neglecting pressure drops in the ideal analysis),

$$\epsilon \equiv \frac{T_5 - T_2}{T_4 - T_2}, \quad 0 \leq \epsilon \leq 1,$$

giving the preheated temperature:

$$T_5 = T_2 + \epsilon (T_4 - T_2).$$

- **Adjusted heat input with regeneration:**

The combustor now heats from  $T_5$  to  $T_3$ ,

$$q_{in, reg} = c_p (T_3 - T_5) = c_p [(T_3 - T_2) - \epsilon (T_4 - T_2)].$$

Compared with the simple cycle heat input  $q_{in} = c_p (T_3 - T_2)$ , regeneration reduces fuel heat by

$$q_{saved} = c_p (T_5 - T_2) = \epsilon c_p (T_4 - T_2).$$

- **Net work and efficiency with regeneration:**

Regeneration does not change the ideal compressor and turbine works, so

$$w_{net, reg} = w_{net} = c_p [(T_3 - T_4) - (T_2 - T_1)].$$

Therefore, the regenerative thermal efficiency is

$$\eta_{th, reg} = \frac{w_{net}}{q_{in, reg}} = \frac{(T_3 - T_4) - (T_2 - T_1)}{(T_3 - T_2) - \epsilon (T_4 - T_2)}.$$

- **Substitution using isentropic relations (optional clarity):**

With  $T_2 = rT_1$  and  $T_4 = \frac{T_3}{r}$ ,

$$\eta_{th, reg} = \frac{T_3(1 - \frac{1}{r}) - T_1(r - 1)}{(T_3 - rT_1) - \varepsilon(\frac{T_3}{r} - rT_1)}.$$

This form shows explicitly how  $\eta_{th, reg}$  rises with  $\varepsilon$ , and how the gain depends on the pressure ratio through  $r$ . Gains are strongest when  $T_4 - T_2$  is large, i.e., at moderate  $r_p$ .

- **Limiting cases and design insight:**

- **No regeneration:**  $\varepsilon = 0 \Rightarrow \eta_{th, reg} = \eta_{th}$ .
- **Perfect regeneration:**  $\varepsilon \rightarrow 1 \Rightarrow q_{in, reg} \rightarrow c_p(T_3 - T_4)$ , so  $\eta_{th, reg} \rightarrow \frac{(T_3 - T_4) - (T_2 - T_1)}{(T_3 - T_4)} = 1 - \frac{T_2 - T_1}{T_3 - T_4}$ .
- **Pressure ratio effect:** As  $r_p$  becomes very large,  $T_4 \rightarrow T_2$ , diminishing the available exhaust-to-compressor temperature difference and reducing the benefit of regeneration.

## Direct answers

- **Simple Brayton net work:**

$$w_{net} = c_p[(T_3 - T_4) - (T_2 - T_1)]$$

- **Simple Brayton efficiency:**

$$\eta_{th} = 1 - \frac{T_4 - T_1}{T_3 - T_2} \quad \text{and} \quad \eta_{th} = 1 - r_p^{-\frac{\gamma-1}{\gamma}}$$

- **Regenerative heat input:**

$$q_{in, reg} = c_p[(T_3 - T_2) - \varepsilon(T_4 - T_2)]$$

- **Regenerative efficiency:**

$$\eta_{th, reg} = \frac{(T_3 - T_4) - (T_2 - T_1)}{(T_3 - T_2) - \varepsilon(T_4 - T_2)} \quad \text{or} \quad \eta_{th, reg} = \frac{T_3(1 - \frac{1}{r}) - T_1(r - 1)}{(T_3 - rT_1) - \varepsilon(\frac{T_3}{r} - rT_1)}.$$

## Notes on reheat and intercooling effects

- **Reheat alone:** Increases  $w_t$  and  $w_{net}$ , but also increases  $q_{in}$ ; efficiency may decrease unless paired with regeneration.
- **Intercooling alone:** Decreases  $w_c$  and increases  $w_{net}$ , but often increases  $q_{in}$ ; efficiency may decrease unless paired with regeneration.
- **Combined strategies:** Intercooling + reheat + regeneration can yield higher  $w_{net}$  and improved  $\eta_{th}$ , balancing added heat with recovered exhaust heat.

## Key takeaways

- **Efficiency (simple Brayton):**  $\eta_{th} = 1 - r_p^{-(\gamma-1)/\gamma}$ .
- **Net work:**  $w_{net} = c_p[(T_3 - T_4) - (T_2 - T_1)]$ .
- **Regeneration:** Reduces  $q_{in}$  via exhaust-heat recovery; stronger gains at moderate  $r_p$ .
- **Reheat:** Raises  $w_{net}$  by increasing average turbine temperature; may lower  $\eta_{th}$  unless combined.
- **Intercooling:** Lowers compressor work; may require regeneration to maintain or improve  $\eta_{th}$ .
- **Optimum  $r_p$  for max  $w_{net}$ :**  $r_p^* = \left(\frac{T_3}{T_1}\right)^{\frac{\gamma}{2(\gamma-1)}}$ .